

# The analyticity region of the hard sphere gas. Improved bounds

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## Abstract

We find an improved estimate of the radius of analyticity of the pressure of the hard-sphere gas in  $d$  dimensions. The estimates are determined by the volume of multidimensional regions that can be numerically computed. For  $d = 2$ , for instance, our estimate is about 40% larger than the classical one.

In a recent paper [4] two of us have shown that it is possible to improve the radius of convergence of the cluster expansion using a tree graph identity due to Penrose [2], see also [5, Section 3]. In this short letter we use the same idea to improve the estimates of the radius of analyticity of the pressure of the hard-sphere gas.

The grand partition function  $\Xi(z, \Lambda)$  of a gas of hard spheres of diameter  $R$  enclosed in a volume  $\Lambda \subset \mathbb{R}$  is given by

$$\Xi(z, \Lambda) = \sum_{n=0}^{\infty} \frac{z^n}{n!} \int_{\Lambda^n} dx_1 \dots dx_n \exp \left\{ - \sum_{1 \leq i < j \leq n} U(x_i - x_j) \right\}$$

with

$$U(x - y) = \begin{cases} 0 & \text{if } |x - y| > R \\ +\infty & \text{if } |x - y| \leq R \end{cases}$$

where  $|x - y|$  denotes the euclidean distance between the sphere centers  $x$  and  $y$ . The corresponding pressure is  $\lim_{\Lambda \rightarrow \mathbb{R}^d} P(z, \Lambda)$  (limit in van Hove sense), where

$$P(z, \Lambda) = \frac{1}{|\Lambda|} \log \Xi(z, \Lambda)$$

( $|\Lambda|$  denotes the volume of the region  $\Lambda$ ). The cluster expansion, in this setting, amounts to writing the preceding logarithm as the power series (see e.g. [1])

$$\log \Xi(z, \Lambda) = \sum_{n=1}^{\infty} \frac{z^n}{n!} \int_{\Lambda^n} dx_1 \dots dx_n \sum_{\substack{g \subset g(x_1, \dots, x_n) \\ g \in G_n}} (-1)^{|g|} \quad (1)$$

where the graph  $g(x_1, \dots, x_n)$  has vertex set  $\{1, \dots, n\}$  and edge set  $E(x_1, \dots, x_n) = \{\{i, j\} : |x_i - x_j| \leq R\}$  (that is, if the spheres centered at  $x_i$  and  $x_j$  intersect),  $G_n$  is the set of all the connected graphs with vertex set  $\{1, \dots, n\}$ , and  $|g|$  denotes the cardinality of the edge set of the graph  $g$ . Only families  $(x_1, \dots, x_n)$  for which  $g(x_1, \dots, x_n)$  is connected contribute to (1); such families represent “clusters” of spheres.

The standard way to estimate the radius of analyticity of the pressure is to obtain a  $\Lambda$ -independent lower bound of the radius of convergence of the series

$$|P|(z, \Lambda) = \frac{1}{|\Lambda|} \sum_{n=1}^{\infty} \frac{|z|^n}{n!} \int_{\Lambda^n} dx_1 \dots dx_n \left| \sum_{\substack{g \subset g(x_1, \dots, x_n) \\ g \in G_n}} (-1)^{|g|} \right|. \quad (2)$$

This strategy leads to the classical estimation (see e.g. [6], Section 4) that the pressure is analytic if

$$|z| < \frac{1}{e V_d(R)}, \quad (3)$$

where  $V_d(R)$  is the volume of the  $d$ -dimensional sphere of radius  $R$  (excluded volume).

Our approach is based on a well known tree identity. Let us denote by  $T_n$  the subset of  $G_n$  formed by all tree graphs with vertex set  $\{1, \dots, n\}$ . Given a tree  $\tau \in T_n$  and a vertex  $i$  of  $\tau$ , we denote by  $d_i$  the *degree* of the vertex  $i$  in  $\tau$ , i.e. the number of edges of  $\tau$  containing  $i$ . We regard the trees  $\tau \in T_n$  as rooted in the vertex 1. This determines the usual partial order of vertices in  $\tau$  by generations: If  $u, v$  are vertices of  $\tau$ , we write  $u \prec v$ —and say that  $u$  precedes  $v$ —if the (unique) path from the root to  $v$  contains  $u$ . If  $\{u, v\}$  is an edge of  $\tau$ , then either  $v \prec u$  or  $u \prec v$ . Let  $\{u, v\}$  be an edge of  $\tau$  and assume without loss of generality that  $u \prec v$ , then  $u$  is called the *predecessor* and  $v$  the *descendant*. Every vertex  $v \in \tau$  has a unique predecessor and  $s_v = d_v - 1$  descendants, except the root that has no predecessor and  $s_v = d_v$  descendants. For each vertex  $v$  of  $\tau$  we denote by  $v'$  the unique predecessor of  $v$  and by  $v^1, \dots, v^{s_v}$  the  $s_v$  descendants of  $v$ . The number  $s_v$  is called the branching factor; vertices with  $s_v = 0$  are called end-points or “leaves”.

Penrose [4] showed that the sum in (1) is equal, up to a sign, to a sum over trees satisfying certain constraints. We shall keep only the “single-vertex” constraints: descendants of a given sphere must be mutually non-intersecting. This implies that

$$\left| \sum_{\substack{g \subset g(x_1, \dots, x_n) \\ g \in G_n}} (-1)^{|g|} \right| \leq \sum_{\tau \in T_n} w_{\tau}(x_1, \dots, x_n) \quad (4)$$

where

$$w_{\tau}(x_1, \dots, x_n) = \begin{cases} 1 & \text{if } |x_v - x_{v'}| \leq R \text{ and } |x_{v^i} - x_{v^j}| > R, \forall v \text{ vertex of } \tau, \{i, j\} \subset \{1, \dots, s_v\}, \\ 0 & \text{otherwise} \end{cases} \quad (5)$$

Hence from (2) and (4) we get

$$|P|(z, \Lambda) \leq \sum_{n=1}^{\infty} \frac{|z|^n}{n!} \sum_{\tau \in T_n} S_{\Lambda}(\tau) \quad (6)$$

with

$$S_{\Lambda}(\tau) = \frac{1}{|\Lambda|} \int_{\Lambda^n} w_{\tau}(x_1, \dots, x_n) dx_1 \dots dx_n. \quad (7)$$

By (5) we have

$$S_\Lambda(\tau) \leq g_d(d_1) \prod_{i=2}^n g_d(d_i - 1) \quad (8)$$

where  $d_i$  is the degree of the vertex  $i$  in  $\tau$ ,

$$g_d(k) = \int_{\substack{|x_i| \leq R \\ |x_i - x_j| > R}} dx_1 \dots dx_k = R^{dk} \int_{\substack{|y_i| \leq 1 \\ |y_i - y_j| > 1}} dy_1 \dots dy_k$$

for  $k$  positive integer, and  $g_d(0) = 1$  by definition. It is convenient to write

$$g_d(k) = [V_d(R)]^k \tilde{g}_d(k) \quad (9)$$

with

$$\tilde{g}_d(k) = \frac{1}{[V_d(1)]^k} \int_{\substack{|y_i| \leq 1 \\ |y_i - y_j| > 1}} dy_1 \dots dy_k \quad (10)$$

for  $k$  positive integer and  $\tilde{g}_d(0) = 1$ . We observe that  $\tilde{g}_d(k) \leq 1$  for all values of  $k$ . From (8)–(10) we conclude that

$$\begin{aligned} S_\Lambda(\tau) &\leq [V_d(R)]^{d_1} \tilde{g}_d(d_1) \prod_{i=2}^n [V_d(R)]^{d_i-1} \tilde{g}_d(d_i - 1) \\ &= [V_d(R)]^{n-1} \tilde{g}_d(d_1) \prod_{i=2}^n \tilde{g}_d(d_i - 1). \end{aligned}$$

The last identity follows from the fact that for every tree of  $n$  vertices,  $d_1 + \dots + d_n = 2n - 2$ . The  $\tau$ -dependence of this last bound is only through the degree of the vertices, hence it leads, upon insertion in (6), to the inequality

$$|P|(z, \Lambda) \leq \frac{1}{V_d(R)} \sum_{n=1}^{\infty} \frac{(|z| V_d(R))^n}{n!} \sum_{\substack{d_1, \dots, d_n \\ d_1 + \dots + d_n = 2n-2}} \tilde{g}_d(d_1) \prod_{i=2}^n \tilde{g}_d(d_i - 1) \frac{(n-2)!}{\prod_{i=1}^n (d_i - 1)!}.$$

The last quotient of factorials is, precisely, the number of trees with  $n$  vertices and fixed degrees  $d_1, \dots, d_n$ , according to Cayley formula.

At this point we can bound the last sum by a power in an obvious manner. The convergence condition so obtained would already be an improvement over the classical estimate (3). We can, however, get an even better result through a trick used by two of us in [3]. We multiply and divide by  $a^{n-1}$  where  $a > 0$  is a parameter to be chosen in an optimal way. This leads us to the inequality

$$\begin{aligned} |P|(z, \Lambda) &\leq \frac{a}{V_d(R)} \sum_{n=1}^{\infty} \frac{(|z| V_d(R))^n}{a^n n(n-1)} \sum_{\substack{d_1, \dots, d_n \\ d_1 + \dots + d_n = 2n-2}} \frac{\tilde{g}_d(d_1) a^{d_1}}{d_1!} \prod_{i=2}^n \frac{\tilde{g}_d(d_i - 1) a^{d_i-1}}{(d_i - 1)!} \\ &\leq \frac{a}{V_d(R)} \sum_{n=1}^{\infty} \frac{1}{n(n-1)} \left( \frac{|z|}{a} [V_d(R)] \left[ \sum_{s \geq 0} \frac{\tilde{g}_d(s) a^s}{s!} \right] \right)^n \end{aligned}$$

The last series converges if

$$|z| V_d(R) \leq \frac{a}{C_d(a)}$$

where

$$C_d(a) = \sum_{s \geq 0} \frac{\tilde{g}_d(s)}{s!} a^s$$

(this is a finite sum!). The pressure is, therefore, analytic if

$$|z| V_d(R) \leq \max_{a>0} \frac{a}{C_d(a)} . \quad (11)$$

This is our new condition.

Let us show that for  $d = 2$  the quantitative improvement given by this condition can be substantial. In this case

$$C_2(a) = \sum_{s=0}^5 \frac{\tilde{g}_2(s)}{s!} a^s$$

where, by definition,  $\tilde{g}_2(0) = \tilde{g}_2(1) = 1$ . The factor  $\tilde{g}_2(2)$  can be explicitly evaluated in terms of straightforward integrals and we get

$$\tilde{g}_2(2) = \frac{3\sqrt{3}}{4\pi}$$

The other terms of the sum can be numerically evaluated using a simple Montecarlo simulation, obtaining

$$\tilde{g}_2(3) = 0,0589 \quad \tilde{g}_2(4) = 0,0013 \quad \tilde{g}_2(5) \leq 0,0001$$

Choosing  $a = \left[ \frac{8\pi}{3\sqrt{3}} \right]^{1/2}$  (a value for which  $\frac{a}{C_d(a)}$  is close to its maximum) we get

$$|z| V_2(R) \leq 0.5107 .$$

This should be compared with the bound  $|z| V_2(R) \leq 0.36787 \dots$  obtained through the classical condition (3).

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